

Solutions

(1) $Z^2 = \frac{|1\rangle\langle 2| |1\rangle\langle 2|}{\langle 2| |1\rangle\langle 2|} = \frac{|1\rangle\langle 2|}{\langle 2| |1\rangle} = Z$ is immediate. The eigenvalues z of Z must therefore obey $z^2 = z$, so that $z=0$ and $z=1$ are the only possibilities. Indeed,

$$Z|1\rangle = \frac{|1\rangle\langle 2| |1\rangle}{\langle 2| |1\rangle} = |1\rangle,$$

$$\langle 2|Z = \frac{\langle 2| |1\rangle\langle 2|}{\langle 2| |1\rangle} = \langle 2|$$

show what is reasonably obvious in the first place, namely that $|1\rangle$ is the eigeket of Z to eigenvalue 1, and $\langle 2|$ is the eigenbra of Z . Every ket $|j\rangle$ with $\langle 2|j\rangle = 0$ is eigeket to eigenvalue 0, and every bra $\langle l|$ with $\langle l|1\rangle = 0$ is eigenbra to eigenvalue 0.

(2) Recall the definitions $\frac{\partial}{\partial x} f = \frac{1}{i\hbar} [f, P]$, $\frac{\partial}{\partial P} f = \frac{1}{i\hbar} [x, f]$ to express the difference as

$$\begin{aligned} \frac{\partial}{\partial x} \frac{\partial}{\partial P} f - \frac{\partial}{\partial P} \frac{\partial}{\partial x} f &= \left(\frac{1}{i\hbar}\right)^2 [[x, f], P] - \left(\frac{1}{i\hbar}\right)^2 [x, [f, P]] \\ &= \left(\frac{1}{i\hbar}\right)^2 \left([[x, f], P] + [[f, P], x] \right) \\ &= \left(\frac{1}{i\hbar}\right)^2 \underbrace{\left([[x, f], P] + [[f, P], x] + [[P, x], f] + [[x, P], f] \right)}_{=0 \text{ as a consequence of the Jacobi identity}} \underbrace{}_{L=0 \text{ because } [xP]=i\hbar \text{ commutes with } f} \\ &= 0. \end{aligned}$$

Conclusion: $\frac{\partial}{\partial x} \frac{\partial}{\partial P} f = \frac{\partial}{\partial P} \frac{\partial}{\partial x} f$; they are the same.

(3) We have $P(t) = P(0)$ and $X(t) = X(0) + \frac{t}{M} P(0)$, so that

$$P(t)^2 = P(0)^2, \quad X(t)^2 = X(0)^2 + \frac{t^2}{M^2} P(0)^2 + \frac{t}{M} (X(0)P(0) + P(0)X(0))$$

and

$$\langle P(t) \rangle = \langle P(0) \rangle = 0,$$

$$\langle P(t)^2 \rangle = \langle P(0)^2 \rangle = p_0^2,$$

$$\langle X(t) \rangle = \langle X(0) \rangle + \frac{t}{M} \langle P(0) \rangle = 0,$$

$$\begin{aligned} \langle X(t)^2 \rangle &= \langle X(0)^2 \rangle + \frac{t^2}{M^2} \langle P(0)^2 \rangle + \frac{t}{M} \langle (X(0)P(0) + P(0)X(0)) \rangle \\ &= x_0^2 + \frac{t^2}{M^2} p_0^2 \end{aligned}$$

follow, so that

$$\delta P(t) = \sqrt{\langle P(t)^2 \rangle - \langle P(t) \rangle^2} = p_0,$$

$$\delta X(t) = \sqrt{\langle X(t)^2 \rangle - \langle X(t) \rangle^2} = \sqrt{x_0^2 + \left(\frac{p_0 t}{M}\right)^2}.$$

The uncertainty relation is, of course, obeyed at all times

$$\delta X(t) \delta P(t) = x_0 p_0 \sqrt{1 + \left(\frac{p_0 t}{M x_0}\right)^2} \geq x_0 p_0 \geq \frac{t}{2}.$$

(4) Write $A^{+n} A^n = f_n(A^\dagger A)$, then $f_0(A^\dagger A) = 1$ and $f_1(A^\dagger A) = A^\dagger A$. Further we have

$$f_n(A^\dagger A) = A^{+n} (A^{+n-1} A^{n-1}) A = A^{+} f_{n-1}(A^\dagger A) A$$

so that we get the recurrence relation

$$f_n(A^+A) = A^+A f_{n-1}(A^+A - 1).$$

After looking at the first few

$$f_2(A^+A) = A^+A(A^+A - 1),$$

$$f_3(A^+A) = A^+A(A^+A - 1)(A^+A - 2)$$

we recognize the pattern and infer that

$$f_n(A^+A) = A^+A(A^+A - 1)(A^+A - 2) \dots (A^+A - n + 1)$$

$$= \prod_{k=0}^{n-1} (A^+A - k) = \frac{(A^+A)!}{(A^+A - n)!}$$

One checks easily that the recurrence relation is obeyed.

(5) We differentiate with respect to the numerical parameter λ and get

$$\frac{\partial}{\partial \lambda} e^{-i\lambda S/t} \vec{R} e^{i\lambda S/t} = e^{-i\lambda S/t} \underbrace{\frac{1}{it} [S, \vec{R}]}_{= -i\hbar \vec{R}} e^{i\lambda S/t},$$

$$= -e^{-i\lambda S/t} \vec{R} e^{i\lambda S/t}.$$

This differential equation is of the simple form $f'(\lambda) = -f(\lambda)$, solved by $f(\lambda) = e^{-\lambda} f(0)$, so that

$$e^{-i\lambda S/t} \vec{R} e^{i\lambda S/t} = e^{-\lambda} \vec{R}.$$

Likewise one gets $e^{-i\lambda S/t} \vec{P} e^{i\lambda S/t} = e^{\lambda} \vec{P}$.

(6) Suppressing any other potential quantum numbers, we have

$$|l\rangle = |l, m\rangle, \quad \vec{L}^2 |l, m\rangle = |l, m\rangle \hbar^2 l(l+1), \quad L_3 |l, m\rangle = |l, m\rangle \hbar m$$

and

$$(L_1 + iL_2) |l, m\rangle \propto |l, m+1\rangle, \quad (\text{or vanishes if } m=l),$$

$$(L_1 + iL_2)^2 |l, m\rangle \propto |l, m+2\rangle \quad (\text{or vanishes if } m \geq l-1).$$

Since $|l, m\rangle$ is orthogonal to $|l, m+1\rangle$ and $|l, m+2\rangle$, we thus have

$$\langle (L_1 + iL_2) \rangle = 0,$$

$$\langle (L_1 + iL_2)^2 \rangle = 0$$

or, after separating the real and imaginary parts,

$$\langle L_1 \rangle = \langle L_2 \rangle = 0,$$

$$\langle (L_1^2 - L_2^2) \rangle = 0, \quad \langle (L_1 L_2 + L_2 L_1) \rangle = 0.$$

With $\langle (L_1^2 + L_2^2) \rangle = \langle (\vec{L}^2 - L_3^2) \rangle = \hbar^2 (l(l+1) - m^2)$

this gives

$$\langle L_1^2 \rangle = \langle L_2^2 \rangle = \frac{\hbar^2}{2} (l(l+1) - m^2)$$

and then

$$\delta L_1 = \delta L_2 = \frac{\hbar}{\sqrt{2}} \sqrt{l(l+1) - m^2}.$$

Since $[L_1, L_2] = i\hbar L_3$, the uncertainty relation for L_1

and L_2 is

$$\delta L_1 \delta L_2 \geq \frac{1}{2} |\langle i [L_1, L_2] \rangle| = \frac{\hbar}{2} |\langle L_3 \rangle|,$$

here:

$$\frac{\hbar^2}{2} (l(l+1) - m^2) \geq \frac{\hbar^2}{2} m,$$

which is tantamount to $l \geq |m|$ and thus obeyed, indeed. //